

ON CERTAIN INTEGRAL TRANSFORM INVOLVING GENERALIZED BESSEL-MAITLAND FUNCTION AND APPLICATIONS

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ABSTRACT. In this article, we establish a new integral formula involving the generalized Bessel-Maitland function defined by Khan et al. [9], which is expressed in terms of generalized (Wright) hypergeometric function. Some interesting and special cases of our main result are also considered.

1. INTRODUCTION

In recent years, many authors (see, e.g., [1-4] have developed numerous integral formulas involving a variety of special functions. Also many integral formulas associated with the Bessel functions of several kinds have been presented (see, e.g., [6-9]). Those integrals involving Bessel-Maitland functions are not only of great interest to the pure mathematics, but they are often of extreme importance in many branches of theoretical and applied physics and engineering (see [20]). Several methods for evaluating infinite or finite integrals involving Bessel-Maitland functions have been known (see, e.g.,[4] and [19]). However, these methods usually work on a case-by-case basis.

Currently, Ghayasuddin and Khan [4], Khan et al. [6-9] gave certain interesting new class of integral formulas involving the generalized Bessel-Maitland function, which are expressed in terms of the generalized (Wright) hypergeometric function. In the present sequel to the aforementioned investigations, we present two generalized integral formulas involving generalized Bessel-Maitland functions, which are expressed in terms of the generalized (Wright) hypergeometric function. Some special cases and the (potential) usefulness of our main results are also considered and remarked, respectively.

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The Bessel-Maitland function $J_\nu^\mu(z)$ [12;Eq.(8.3)] defined by the following series representation:

$$J_\nu^\mu(z) = \sum_{n=0}^{\infty} \frac{(-z)^n}{n! \Gamma(\mu n + \nu + 1)} = \phi(\mu, \nu + 1; -z). \quad (1.1)$$

Singh et al. [19] introduced the following generalization of Bessel-Maitland function as:

$$J_{\nu,q}^{\mu,\gamma} = \sum_{n=0}^{\infty} \frac{(\gamma)_{qn} (-z)^n}{\Gamma(\mu n + \nu + 1) n!}, \quad (1.2)$$

where $\mu, \nu, \gamma \in \mathbb{C}$, $\operatorname{Re}(\mu) \geq 0$, $\operatorname{Re}(\nu) \geq -1$, $\operatorname{Re}(\gamma) \geq 0$ and $q \in (0, 1) \cup \mathbb{N}$ and $(\gamma)_0 = 1$, $(\gamma)_{qn} = \frac{\Gamma(\gamma+qn)}{\Gamma(\gamma)}$ denotes the generalized Pochhammer symbol.

Recently, Ghayasuddin and Khan [4] introduced and investigated generalized Bessel-Maitland function defined as

$$J_{\nu,\gamma,\delta}^{\mu,q,p}(z) = \sum_{n=0}^{\infty} \frac{(\gamma)_{qn} (-z)^n}{\Gamma(\mu n + \nu + 1) (\delta)_{pn}}, \quad (1.3)$$

where $\mu, \nu, \gamma, \delta \in \mathbb{C}$, $\operatorname{Re}(\mu) \geq 0$, $\operatorname{Re}(\nu) \geq -1$, $\operatorname{Re}(\gamma) \geq 0$, $\operatorname{Re}(\delta) \geq 0$; $p, q > 0$ and $q < \operatorname{Re}(\alpha) + p$.

In particular Khan et al. [9] introduced and investigated a new extension of Bessel-Maitland function as follows:

$$J_{\alpha,\beta,\nu,\sigma,\delta,p}^{\mu,\rho,\gamma,q}(z) = \sum_{n=0}^{\infty} \frac{(\mu)_{pn} (\gamma)_{qn} (-z)^n}{\Gamma(n\beta + \alpha + 1) (\delta)_{pn} (\nu)_{n\sigma}}, \quad (1.4)$$

where $\alpha, \beta, \mu, \rho, \nu, \gamma, \sigma, \delta \in \mathbb{C}$; $\operatorname{Re}(\beta) > 0$, $\operatorname{Re}(\rho) > 0$, $\operatorname{Re}(\mu) > 0$, $\operatorname{Re}(\nu) > 0$, $\operatorname{Re}(\alpha) \geq -1$, $\operatorname{Re}(\gamma) > 0$, $\operatorname{Re}(\delta) > 0$, $\operatorname{Re}(\sigma) > 0$; $p, q > 0$, and $q < \operatorname{Re}(\alpha) + p$.

Relation with Mittag-Leffler functions:

(i) On replacing α by $\alpha - 1$ in (1.4), we get the following interesting relation:

$$J_{\alpha-1,\beta,\nu,\sigma,\delta,p}^{\mu,\rho,\gamma,q}(-z) = E_{\alpha,\beta,\nu,\sigma,\delta,p}^{\mu,\rho,\gamma,q}(z), \quad (1.5)$$

where $E_{\alpha,\beta,\nu,\sigma,\delta,p}^{\mu,\rho,\gamma,q}(z)$ is the Mittag-Leffler function defined by Khan and Ahmed [10].

(ii) On setting $\mu = \nu = \sigma = \rho = 1$ and replacing α by $\alpha - 1$ in (1.4), we get

$$J_{\alpha-1,\beta,1,1,\delta,p}^{1,1,\gamma,q}(-z) = E_{\alpha,\beta,p}^{\gamma,\delta,q}(z), \quad (1.6)$$

where $E_{\alpha,\beta,p}^{\gamma,\delta,q}(z)$ is the Mittag-Leffler function defined by Salim and Faraz [17].

(iii) On setting $\mu = \nu = \sigma = \rho = \delta = p = 1$ and replacing α by $\alpha - 1$ in (1.4), we get

$$J_{\alpha-1,\beta,1,1,1,1}^{1,1,\gamma,q}(-z) = E_{\alpha,\beta}^{\gamma,q}(z), \quad (1.7)$$

where $E_{\alpha,\beta}^{\gamma,q}(z)$ is the Mittag-Leffler function defined by Shukla and Prajapati [18].

(iv) On setting $\mu = \nu = \sigma = \rho = \delta = p = q = 1$ and replacing α by $\alpha - 1$ in (1.4), we get

$$J_{\alpha-1,\beta,1,1,1}^{1,1,\gamma,1}(-z) = E_{\alpha,\beta}^{\gamma}(z), \quad (1.8)$$

where $E_{\alpha,\beta}^{\gamma}(z)$ is the Mittag-Leffler function defined by Prabhakar [14].

(v) On setting $\mu = \nu = \sigma = \rho = \delta = \gamma = p = q = 1$ and replacing α by $\alpha - 1$ in (1.4), we get

$$J_{\alpha-1,\beta,1,1,1}^{1,1,1,1}(-z) = E_{\alpha,\beta}(z), \quad (1.9)$$

where $E_{\alpha,\beta}(z)$ is the Mittag-Leffler function defined by Wiman [21].

(vi) On setting $\mu = \nu = \sigma = \rho = \delta = \gamma = p = q = 1$, $\alpha = 0$ and replacing α by $\alpha - 1$ in (1.4), we get

$$J_{0,\beta,1,1,1}^{1,1,1,1}(-z) = E_{\beta}(z), \quad (1.10)$$

where $E_{\beta}(z)$ is the Mittag-Leffler function defined by Ghosta Mittag-Leffler [11].

The generalization of the generalized hypergeometric series ${}_pF_q$ is due to Fox [5] and Wright ([21], [22], [23]) who studied the asymptotic expansion of the generalized (Wright) hypergeometric function defined by (see [17, p.21]; see also [15]):

$${}_p\Psi_q \left[\begin{array}{l} (\alpha_1, A_1), \dots, (\alpha_p, A_p); \\ (\beta_1, B_1), \dots, (\beta_q, B_q); \end{array} z \right] = \sum_{k=0}^{\infty} \frac{\prod_{j=1}^p \Gamma(\alpha_j + A_j k)}{\prod_{j=1}^q \Gamma(\beta_j + B_j k)} \frac{z^k}{k!}, \quad (1.11)$$

where the coefficients A_1, \dots, A_p and B_1, \dots, B_q are positive real numbers such as

$$(i) 1 + \sum_{j=1}^q B_j - \sum_{j=1}^p A_j > 0 \text{ and } 0 < |z| < \infty; z \neq 0. \quad (1.12)$$

$$(ii) 1 + \sum_{j=1}^q B_j - \sum_{j=1}^p A_j = 0 \text{ and } 0 < |z| < A_1^{-A_1} \dots A_p^{-A_p} B_1^{B_1} \dots B_q^{B_q}. \quad (1.13)$$

A special case of (1.11) is

$${}_p\Psi_q \left[\begin{array}{l} (\alpha_1, 1), \dots, (\alpha_p, 1); \\ (\beta_1, 1), \dots, (\beta_q, 1); \end{array} z \right] = \frac{\prod_{j=1}^p \Gamma(\alpha_j)}{\prod_{j=1}^q \Gamma(\beta_j)} {}_pF_q \left[\begin{array}{l} \alpha_1, \dots, \alpha_p; \\ \beta_1, \dots, \beta_q; \end{array} z \right], \quad (1.14)$$

where ${}_pF_q$ is the generalized hypergeometric series defined by [15]

$$\begin{aligned} {}_pF_q \left[\begin{array}{l} \alpha_1, \dots, \alpha_p; \\ \beta_1, \dots, \beta_q; \end{array} z \right] &= \sum_{n=0}^{\infty} \frac{(\alpha_1)_n \cdots (\alpha_p)_n}{(\beta_1)_n \cdots (\beta_q)_n} \frac{z^n}{n!} \\ &= {}_pF_q(\alpha_1, \dots, \alpha_p; \beta_1, \dots, \beta_q; z), \end{aligned} \quad (1.15)$$

where $(\lambda)_n$ is the Pochhammer's symbol (see [15]).

For our present investigation the following interesting and useful result due to Obhettinger [13] will be required:

$$\int_0^\infty x^{\mu-1} (x + a + \sqrt{x^2 + 2ax})^{-\lambda} dx = 2\lambda a^{-\lambda} \left(\frac{a}{2}\right)^\mu \frac{\Gamma(2\mu)\Gamma(\lambda - \mu)}{\Gamma(1 + \lambda + \mu)}, \quad (1.16)$$

provided $0 < \Re(\mu) < \Re(\lambda)$.

2. MAIN RESULT

Theorem 2.1. If $\alpha, \beta, \gamma, \mu, \eta, \nu, \sigma, \rho, \delta, \lambda \in \mathbb{C}$, $\sigma + \eta + p - \rho - q > 0$, $\Re(\alpha) > 0$, $\Re(\beta) > 0$, $\Re(\gamma) > 0$, $\Re(\mu) > 0$, $\Re(\lambda) \geq -1$, $\Re(\nu) > 0$, $\Re(\delta) > 0$, $\Re(\sigma) > 0$, $\Re(\rho) > 0$, $\Re(\eta) > 0$, $p, q > 0$ and $q < \Re(\alpha) + p$, then

$$\begin{aligned} & \int_0^\infty x^{\alpha-1} (x + a + \sqrt{x^2 + 2ax})^{-\beta} J_{\eta, \lambda, \nu, \sigma, \delta, p}^{\mu, \rho, \gamma, q} \left(\frac{y}{x + a + \sqrt{x^2 + 2ax}} \right) dx \\ &= 2^{1-\alpha} a^{\alpha-\beta} \frac{\Gamma(2\alpha)\Gamma(\nu)\Gamma(\delta)}{\Gamma(\mu)\Gamma(\gamma)} {}_5\Psi_5 \left[\begin{array}{c} (\mu, \rho), (\gamma, q), (\beta - \alpha, 1), (\beta + 1, 1), (1, 1); \\ (\nu, \sigma), (\lambda + 1, \eta), (\beta, 1), (1 + \beta + \alpha, 1), (\delta, p); \end{array} \middle| \frac{-y}{a} \right]. \end{aligned} \quad (2.1)$$

Proof. In order to establish our main integral (2.1), we denote the left hand side of (2.1) by I and then by using (1.4), to get:

$$\begin{aligned} I &= \int_0^\infty x^{\alpha-1} (x + a + \sqrt{x^2 + 2ax})^{-\beta} \\ &\times \sum_{n=0}^{\infty} \frac{(\mu)_{pn}(\gamma)_{qn}}{\Gamma(\eta n + \lambda + 1)(\nu)_{\sigma n}(\delta)_{pn}} \left(\frac{-y}{x + a + \sqrt{x^2 + 2ax}} \right)^n dx. \end{aligned} \quad (2.2)$$

Evaluating the above integral with the help of (1.16) and after little simplification, we found:

$$\begin{aligned} I &= 2^{1-\alpha} a^{\alpha-\beta} \frac{\Gamma(2\alpha)\Gamma(\nu)\Gamma(\delta)}{\Gamma(\mu)\Gamma(\gamma)} \\ &\times \sum_{n=0}^{\infty} \frac{\Gamma(\mu + \rho n)\Gamma(\gamma + q n)\Gamma(\beta + n - \alpha)\Gamma(\beta + n + 1)\Gamma(n + 1)}{\Gamma(\nu + \sigma n)\Gamma(\eta n + \lambda + 1)\Gamma(\beta + n)\Gamma(1 + \beta + n + \alpha)\Gamma(\delta + pn)n!} \left(\frac{-y}{a} \right)^n. \end{aligned} \quad (2.3)$$

Which upon using (1.11) yields (2.1). This completes the proof of our main result.

Variation of (2.1): If the conditions of our main result be satisfied, then the following integral formula holds true:

$$\begin{aligned} & \int_0^\infty x^{\alpha-1} (x + a + \sqrt{x^2 + 2ax})^{-\beta} J_{\eta, \lambda, \nu, \sigma, \delta, p}^{\mu, \rho, \gamma, q} \left(\frac{y}{x + a + \sqrt{x^2 + 2ax}} \right) dx \\ &= 2^{1-\alpha} a^{\alpha-\beta} \frac{\Gamma(2\alpha)\Gamma(\beta + 1)\Gamma(\beta - \alpha)}{\Gamma(\beta)\Gamma(\lambda + 1)\Gamma(1 + \beta + \alpha)} \\ &\times {}_{\rho+q+3}F_{\sigma+\eta+p+2} \left[\begin{array}{c} \Delta(\rho, \mu), \Delta(q, \gamma), \beta + 1, \beta - \alpha, 1; \\ \Delta(\sigma, \nu), \Delta(p; \delta), \Delta(\eta, \lambda + 1), \beta, 1 + \beta + \alpha; \end{array} \middle| \frac{-\rho^\rho q^q y}{\sigma^\sigma \eta^\eta p^p a} \right], \end{aligned} \quad (2.4)$$

where $\Delta(m; l)$ abbreviates the array of m parameters $\frac{l}{m}, \frac{l+m}{m}, \dots, \frac{l+m-1}{m}$, $m \geq 1$.

Proof: By writing the right hand side of (2.4) in the original summation and using the result

$$\Gamma(\alpha + n) = \Gamma(\alpha)(\alpha)_n$$

and

$$(l)_{kn} = k^{kn} \left(\frac{l}{k} \right)_n \left(\frac{l+1}{k} \right)_n \dots \left(\frac{l+k-1}{k} \right)_n,$$

(Gauss multiplication theorem) in (2.3) and assuming up the given series with the help of (1.15), we easily arrive at our required result (2.4).

3. SPECIAL CASES

(i). On replacing λ by $\lambda - 1$ in (2.1) and then by using (1.5), we obtain:

$$\begin{aligned} & \int_0^\infty x^{\alpha-1} (x + a + \sqrt{x^2 + 2ax})^{-\beta} E_{\eta, \lambda, \nu, \sigma, \delta, p}^{\mu, \rho, \gamma, q} \left(\frac{y}{x + a + \sqrt{x^2 + 2ax}} \right) dx \\ &= 2^{1-\alpha} a^{\alpha-\beta} \frac{\Gamma(2\alpha)\Gamma(\nu)\Gamma(\delta)}{\Gamma(\mu)\Gamma(\gamma)} {}_5\Psi_5 \left[\begin{matrix} (\mu, \rho), (\gamma, q), (\beta - \alpha, 1), (\beta + 1, 1), (1, 1); \\ (\nu, \sigma), (\lambda, \eta), (\beta, 1), (1 + \beta + \alpha, 1), (\delta, p); \end{matrix} \frac{y}{a} \right], \end{aligned} \quad (3.1)$$

where $\alpha, \beta, \gamma, \mu, \eta, \nu, \sigma, \rho, \delta, \lambda \in \mathbb{C}$, $\Re(\alpha) > 0, \Re(\beta) > 0, \Re(\gamma) > 0, \Re(\mu) > 0, \Re(\nu) > 0, \Re(\lambda) > 0, \Re(\delta) > 0, \Re(\sigma) > 0, \Re(\rho) > 0, \Re(\eta) > 0$, $p, q > 0$ and $q < \Re(\alpha) + p$.

(ii). On replacing λ by $\lambda - 1$ in (2.4) and then by using (1.5), we attain:

$$\begin{aligned} & \int_0^\infty x^{\alpha-1} (x + a + \sqrt{x^2 + 2ax})^{-\beta} E_{\eta, \lambda, \nu, \sigma, \delta, p}^{\mu, \rho, \gamma, q} \left(\frac{y}{x + a + \sqrt{x^2 + 2ax}} \right) dx \\ &= 2^{1-\alpha} a^{\alpha-\beta} \frac{\Gamma(2\alpha)\Gamma(\beta + 1)\Gamma(\beta - \alpha)}{\Gamma(\beta)\Gamma(\lambda)\Gamma(1 + \beta + \alpha)} \\ & \times {}_{\rho+q+3}F_{\sigma+\eta+p+2} \left[\begin{matrix} \Delta(\rho, \mu), \Delta(q, \gamma), \beta + 1, \beta - \alpha, 1; \\ \Delta(\sigma, \nu), \Delta(p, \delta), \Delta(\eta, \lambda), \beta, 1 + \beta + \alpha; \end{matrix} \frac{\rho^\rho q^q y}{\sigma^\sigma \eta^\eta p^p a} \right], \end{aligned} \quad (3.2)$$

where $\alpha, \beta, \gamma, \mu, \eta, \nu, \sigma, \rho, \delta, \lambda \in \mathbb{C}$, $\Re(\alpha) > 0, \Re(\beta) > 0, \Re(\gamma) > 0, \Re(\mu) > 0, \Re(\nu) > 0, \Re(\lambda) > 0, \Re(\delta) > 0, \Re(\sigma) > 0, \Re(\rho) > 0, \Re(\eta) > 0$, $p, q > 0$ and $q < \Re(\alpha) + p$.

(iii). On setting $\mu = \nu = \rho = \sigma = 1$ and replacing λ by $\lambda - 1$ in (2.1) and then by using (1.6), we find:

$$\begin{aligned} & \int_0^\infty x^{\alpha-1} (x+a+\sqrt{x^2+2ax})^{-\beta} E_{\eta,\lambda,p}^{\gamma,\delta,q} \left(\frac{y}{x+a+\sqrt{x^2+2ax}} \right) dx \\ &= 2^{1-\alpha} a^{\alpha-\beta} \frac{\Gamma(2\alpha)\Gamma(\delta)}{\Gamma(\gamma)} {}_4\Psi_4 \left[\begin{matrix} (\gamma, q), (\beta-\alpha, 1), (\beta+1, 1), (1, 1); \\ (\lambda, \eta), (\beta, 1), (1+\beta+\alpha, 1), (\delta, p); \end{matrix} \frac{y}{a} \right], \quad (3.3) \end{aligned}$$

where $\alpha, \beta, \gamma, \eta, \delta, \lambda \in \mathbb{C}$, $\Re(\alpha) > 0, \Re(\beta) > 0, \Re(\gamma) > 0, \Re(\lambda) > 0, \Re(\delta) > 0, \Re(\eta) > 0$, $p, q > 0$ and $q < \Re(\alpha) + p$.

(iv). On setting $\mu = \nu = \rho = \sigma = 1$ and replacing λ by $\lambda - 1$ in (2.4) and then by using (1.6), we get:

$$\begin{aligned} & \int_0^\infty x^{\alpha-1} (x+a+\sqrt{x^2+2ax})^{-\beta} E_{\eta,\lambda,p}^{\gamma,\delta,q} \left(\frac{y}{x+a+\sqrt{x^2+2ax}} \right) dx \\ &= 2^{1-\alpha} a^{\alpha-\beta} \frac{\Gamma(2\alpha)\Gamma(\beta+1)\Gamma(\beta-\alpha)}{\Gamma(\beta)\Gamma(\lambda)\Gamma(1+\beta+\alpha)} \\ & \times {}_{q+3}F_{\eta+p+2} \left[\begin{matrix} \Delta(q, \gamma), \beta+1, \beta-\alpha, 1; \\ \Delta(p, \delta), \Delta(\eta, \lambda), \beta, 1+\beta+\alpha; \end{matrix} \frac{q^q y}{\eta^\eta p^p a} \right], \quad (3.4) \end{aligned}$$

where $\alpha, \beta, \gamma, \eta, \delta, \lambda \in \mathbb{C}$, $\Re(\alpha) > 0, \Re(\beta) > 0, \Re(\gamma) > 0, \Re(\lambda) > 0, \Re(\delta) > 0, \Re(\eta) > 0$, $p, q > 0$ and $q < \Re(\alpha) + p$.

(v). On setting $\mu = \nu = \rho = \sigma = \delta = p = 1$ and replacing λ by $\lambda - 1$ in (2.1) and then by using (1.7), we find:

$$\begin{aligned} & \int_0^\infty x^{\alpha-1} (x+a+\sqrt{x^2+2ax})^{-\beta} E_{\eta,\lambda}^{\gamma,q} \left(\frac{y}{x+a+\sqrt{x^2+2ax}} \right) dx \\ &= 2^{1-\alpha} a^{\alpha-\beta} \frac{\Gamma(2\alpha)}{\Gamma(\gamma)} {}_4\Psi_3 \left[\begin{matrix} (\gamma, q), (\beta-\alpha, 1), (\beta+1, 1), (1, 1); \\ (\beta, 1), (1+\beta+\alpha, 1), (\lambda, \eta); \end{matrix} \frac{y}{a} \right], \quad (3.5) \end{aligned}$$

where $\alpha, \beta, \gamma, \eta, \lambda \in \mathbb{C}$, $\Re(\alpha) > 0, \Re(\beta) > 0, \Re(\gamma) > 0, \Re(\lambda) > 0, \Re(\eta) > 0$, $q \in (0, 1) \cup \mathbb{N}$.

(vi). On setting $\mu = \nu = \rho = \sigma = \delta = p = 1$ and replacing λ by $\lambda - 1$ in (2.4) and then by using (1.7), we acquire:

$$\begin{aligned} & \int_0^\infty x^{\alpha-1} (x+a+\sqrt{x^2+2ax})^{-\beta} E_{\eta,\lambda}^{\gamma,q} \left(\frac{y}{x+a+\sqrt{x^2+2ax}} \right) dx \\ &= 2^{1-\alpha} a^{\alpha-\beta} \frac{\Gamma(2\alpha)\Gamma(\beta+1)\Gamma(\beta-\alpha)}{\Gamma(\beta)\Gamma(\lambda)\Gamma(1+\beta+\alpha)} \\ & \times {}_{q+3}F_{\eta+2} \left[\begin{matrix} \Delta(q; \gamma), \beta+1, \beta-\alpha, 1; \\ \Delta(\eta, \lambda), \beta, 1+\beta+\alpha; \end{matrix} \frac{q^q y}{\eta^\eta a} \right], \quad (3.6) \end{aligned}$$

where $\alpha, \beta, \gamma, \eta, \lambda \in \mathbb{C}$, $\Re(\alpha) > 0, \Re(\beta) > 0, \Re(\gamma) > 0, \Re(\lambda) > 0, \Re(\eta) > 0$, $q \in (0, 1) \cup \mathbb{N}$.

(vii). On setting $\mu = \nu = \rho = \sigma = \delta = p = q = 1$ and replacing λ by $\lambda - 1$ in (2.1) and then by using (1.8), we obtain:

$$\begin{aligned} & \int_0^\infty x^{\alpha-1} (x + a + \sqrt{x^2 + 2ax})^{-\beta} E_{\eta, \lambda}^\gamma \left(\frac{y}{x + a + \sqrt{x^2 + 2ax}} \right) dx \\ &= 2^{1-\alpha} a^{\alpha-\beta} \frac{\Gamma(2\alpha)}{\Gamma(\gamma)} {}_4\Psi_3 \left[\begin{matrix} (\gamma, 1), (\beta - \alpha, 1), (\beta + 1, 1), (1, 1); \\ (\beta, 1), (1 + \beta + \alpha, 1), (\lambda, \eta); \end{matrix} \frac{y}{a} \right], \end{aligned} \quad (3.7)$$

where $\alpha, \beta, \gamma, \eta, \lambda \in \mathbb{C}$, $\Re(\alpha) > 0, \Re(\beta) > 0, \Re(\gamma) > 0, \Re(\lambda) > 0, \Re(\eta) > 0$.

(viii). On setting $\mu = \nu = \rho = \sigma = \delta = p = q = 1$ and replacing ν by $\nu - 1$ in (2.1) and then by using (1.7), we get:

$$\begin{aligned} & \int_0^\infty x^{\alpha-1} (x + a + \sqrt{x^2 + 2ax})^{-\beta} E_{\eta, \lambda}^\gamma \left(\frac{y}{x + a + \sqrt{x^2 + 2ax}} \right) dx \\ &= 2^{1-\alpha} a^{\alpha-\beta} \frac{\Gamma(2\alpha)\Gamma(\beta+1)\Gamma(\beta-\alpha)}{\Gamma(\beta)\Gamma(\lambda)\Gamma(1+\beta+\alpha)} \\ & \times {}_4F_{\eta+2} \left[\begin{matrix} \gamma, \beta + 1, \beta - \alpha, 1; \\ \Delta(\eta, \lambda), 1 + \beta + \alpha, \beta; \end{matrix} \frac{y}{\eta^\eta a} \right], \end{aligned} \quad (3.8)$$

where $\alpha, \beta, \gamma, \eta, \lambda \in \mathbb{C}$, $\Re(\alpha) > 0, \Re(\beta) > 0, \Re(\gamma) > 0, \Re(\lambda) > 0, \Re(\eta) > 0$.

(ix). On setting $\mu = \nu = \rho = \eta = p = q = \delta = \gamma = 1$ in (2.1) and then by using (1.8), we attain:

$$\begin{aligned} & \int_0^\infty x^{\alpha-1} (x + a + \sqrt{x^2 + 2ax})^{-\beta} E_{\eta, \lambda}^\gamma \left(\frac{y}{x + a + \sqrt{x^2 + 2ax}} \right) dx \\ &= 2^{1-\alpha} a^{\alpha-\beta} \Gamma(2\alpha) {}_3\Psi_3 \left[\begin{matrix} (\beta - \alpha, 1), (\beta + 1, 1), (1, 1); \\ (\beta, 1), (1 + \beta + \alpha, 1), (\lambda, \eta); \end{matrix} \frac{y}{a} \right], \end{aligned} \quad (3.9)$$

where $\alpha, \beta, \eta, \lambda \in \mathbb{C}$, $\Re(\alpha) > 0, \Re(\beta) > 0, \Re(\lambda) > 0, \Re(\eta) > 0$.

(x). On setting $\mu = \nu = \rho = \eta = p = q = \delta = \gamma = 1$ in (2.1) and then by using (1.8), we obtain:

$$\begin{aligned} & \int_0^\infty x^{\alpha-1} (x + a + \sqrt{x^2 + 2ax})^{-\beta} E_{\eta, \lambda}^\gamma \left(\frac{y}{x + a + \sqrt{x^2 + 2ax}} \right) dx \\ &= 2^{1-\alpha} a^{\alpha-\beta} \frac{\Gamma(2\alpha)\Gamma(\beta+1)\Gamma(\beta-\alpha)}{\Gamma(\beta)\Gamma(\lambda)\Gamma(1+\beta+\alpha)} \end{aligned}$$

$$\times {}_3F_{\eta+2} \left[\begin{array}{c} \beta+1, \beta-\alpha, 1; \\ \Delta(\eta, \lambda), 1+\beta+\alpha, \beta; \end{array} \frac{y}{\eta^\eta a} \right], \quad (3.10)$$

where $\alpha, \beta, \eta, \lambda \in \mathbb{C}$, $\Re(\alpha) > 0$, $\Re(\beta) > 0$, $\Re(\lambda) > 0$, $\Re(\eta) > 0$.

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