

A NEW CLASS OF INTEGRALS INVOLVING EXTENDED MITTAG-LEFFLER FUNCTIONS

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ABSTRACT. The main aim of this paper is to establish two generalized integral formulas involving the extended Mittag-Leffler function based on the well known Lavoie and Trottier integral formula and the results are express in term of extended Wright-type function. Also, we establish certain special cases of our main result.

1. INTRODUCTION

The Mittag-Leffler functions $E_\alpha(z)$ and $E_{\alpha,\beta}(z)$ are defined by the following series:

$$E_\alpha(z) = \sum_{n=0}^{\infty} \frac{z^n}{\Gamma(\alpha n + 1)}, \quad z \in \mathbb{C}, \Re(\alpha) > 0, \quad (1)$$

and

$$E_{\alpha,\beta}(z) = \sum_{n=0}^{\infty} \frac{z^n}{\Gamma(\alpha n + \beta)}, \quad z, \beta \in \mathbb{C}, \Re(\alpha) > 0, \quad (2)$$

respectively. For the generalizations and applications of Mittag-Leffler function, the readers can refer to the recent work of researchers ([5, 6, 7, 8, 9, 11, 23, 25, 33]), Kilbas et al. ([12], Chapter 1) and Saigo and Kilbas [13]. In recent years, the Mittag-Leffler function and some of its variety generalizations have been numerically established in the complex plane ([10, 27]). A new generalization of the Mittag-Leffler functions $E_{\alpha,\beta}(z)$ of (2) have defined by Prabhakar [24] as follows:

$$E_{\alpha,\beta}^\gamma(z) = \sum_{n=0}^{\infty} \frac{(\gamma)_n}{\Gamma(\alpha n + \beta)} \frac{z^n}{n!}, \quad z, \beta \in \mathbb{C}, \Re(\alpha) > 0, \quad (3)$$

where $(\gamma)_n$ denote the well known Pochhammer Symbol which is defined by:

$$(\gamma)_n = \begin{cases} 1, & (n = 0, \gamma \in \mathbb{C}), \\ \gamma(\gamma + 1)(\gamma + 2)\dots(\gamma + n - 1) & (n \in \mathbb{N}, \gamma \in \mathbb{C}). \end{cases}$$

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In fact, the following special cases are satisfied:

$$E_{\alpha,\beta}^1(z) = E_{\alpha,\beta}(z), E_{\alpha,1}^1(z) = E_{\alpha,\beta}(z). \quad (4)$$

Recently, many researchers have established the significance and great consideration of Mittag-Leffler functions in the theory of special functions for exploring the generalizations and some applications. Various extensions for these functions are found in ([9],[28, 29, 30]). Srivastava and Tomovski [32] have defined further generalization of the Mittag-Leffler function $E_{\alpha,\beta}^1(z)$ of (3), which is defined as:

$$E_{\alpha,\beta}^{\gamma,k}(z) = \sum_{n=0}^{\infty} \frac{(\gamma)_{nk}}{\Gamma(\alpha n + \beta)} \frac{z^n}{n!}, \quad (5)$$

where $z, \alpha, \beta \in \mathbb{C}; \Re(\alpha) > \max\{0, \Re(k) - 1\}; \Re(k) > 0$. Very recently, Ozarslan and Yilmaz [22] have investigated an extended Mittag-Leffler function $E_{\alpha,\beta}^{\gamma,c}(z : p)$ which is defined as:

$$E_{\alpha,\beta}^{\gamma;c}(z) = \sum_{n=0}^{\infty} \frac{\mathbf{B}_p(\gamma + n, c - \gamma)(c)_n}{\mathbf{B}(\gamma, c - \gamma)\Gamma(\alpha n + \beta)} \frac{z^n}{n!}, \quad (6)$$

where $p > 0, \Re(c) > \Re(\gamma) > 0$ and $B_p(x, y)$ is extended beta function defined in Mainardi as follows:

$$B_p(x, y) = \int_0^t t^{(x-1)}(1-t)^{(y-1)} e^{-(p/t(1-t))} dt, \quad (7)$$

where $\Re(p) > 0, \Re(x) > 0$ and $\Re(y) > 0$. If $p = 0$, then $B_p(x, y)$ reduces to the following beta function:

$$B(x, y) = \int_0^t t^{(x-1)}(1-t)^{(y-1)} dt. \quad (8)$$

The notion of the extended beta functions given by Chaudhry et al. [2] and further analyzed in [16, 18] The gamma function is defined by:

$$\Gamma(z) = \int_0^t t^{(z-1)} e^{(-t)} dt; \quad \Re(z) > 0. \quad (9)$$

By inspection, we conclude the following relation

$$\Gamma(z+1) = z\Gamma(z). \quad (10)$$

The generalized hypergeometric function ${}_pF_q(z)$ is defined in [6] as:

$$\begin{aligned} {}_pF_q(z) &= {}_pF_q \left[\begin{array}{l} (a_1), (a_2), \dots, (a_p); \\ (b_1), (b_2), \dots, (b_q); \end{array} z, p \right] \\ &= \sum_{n=0}^{\infty} \frac{(a_1)_n(a_2)_n \dots (a_q)_n z^n}{(b_1)_n(b_2)_n \dots (b_p)_n n!} \end{aligned} \quad (11)$$

where $\alpha_i, \beta_i \in \mathbb{C}; i = 1, 2, \dots, p; j = 1, 2, \dots, q$ and $b_j \neq 0, -1, -2, \dots$ and $(z)_n$ is the Pochhammer symbols.

Sharma and Devi [26] defined extended Wright type function in the following series:

$$\begin{aligned} {}_{p+1}F_{q+1}(z; p) &= {}_{p+1}F_{q+1} \left[\begin{array}{c} (a_i, A_i)_{1,p}, (\gamma, 1); \\ (b_j, \beta_j)_{1,p}, (c, 1); \end{array} z \right] \\ &= \sum_{n=0}^{\infty} \frac{\Gamma(a_1 + A_1 n) \dots \Gamma(a_p + A_p n) \mathbf{B}_p(\gamma + n, c - \gamma) z^n}{\Gamma(b_1 + \beta_1 n) \dots \Gamma(b_q + \beta_q n) \Gamma(c - \gamma) n!} \quad (12) \end{aligned}$$

where $i = 1, 2, \dots, p, j = 1, 2, \dots, q, \Re(p) > 0, \Re(c) > \Re(\gamma) > 0$ and $p \geq 0$. Recently, Mittal et al. [17] defined an extended generalized Mittag-Leffler function as:

$$E_{\alpha, \beta}^{\gamma, q; c}(z; p) = \sum_{n=0}^{\infty} \frac{\mathbf{B}_p(\gamma + nq, c - \gamma)(c)_{nq}}{\mathbf{B}(\gamma, c - \gamma)\Gamma(\alpha n + \beta)} \frac{z^n}{n!}, \quad (13)$$

where $\alpha, \beta, \gamma \in \mathbb{C}, \Re(\alpha) > 0, \Re(\beta) > 0, \Re(\gamma) > 0, q > 0$ and $B_p(x, y)$ is an extended beta function defined in (7). In this paper, we derive new integral formulas involving the Mittag-Leffler function (6) and (13). For the present investigation, we need the following result of Lavoie and Trottier [14]

$$\int_0^1 x^{(a-1)} (1-x)^{(2b-1)} \left(1 - \frac{x}{3}\right)^{(2a-1)} \left(1 - \frac{x}{4}\right)^{(b-1)} dx = \left(\frac{2}{3}\right)^{2a} \frac{\Gamma(a)\Gamma(b)}{\Gamma(a+b)}, \quad (14)$$

where $\Re(a) > 0, \Re(b) > 0$.

For more details about the recent works in the field of dynamical systems theory, stochastic systems, non-equilibrium statistical mechanics and quantum mechanics, the readers may refer to the recent work of the researchers [1, 3, 4, 15, 19, 20, 21] and the references cited therein.

2. MAIN RESULT

In this section, the generalized integral formulas involving the Mittag-Leffler function (6) and (13), are established here by inserting with the suitable argument in the integrand of (14).

Theorem 1. Let $y, c, \rho, j, \alpha, \beta, \gamma \in \mathbb{C}$; with $\Re(\alpha) > \max\{0, \Re(c) - 1\}; \Re(c) > 0, \Re(\gamma) > 0, \Re(2\rho + j) > 0$ and $x > 0$. Then the following formula holds true:

$$\begin{aligned} &\int_0^1 x^{(\rho+j-1)} (1-x)^{(2\rho-1)} \left(1 - \frac{x}{3}\right)^{(2(\rho+j)-1)} \left(1 - \frac{x}{4}\right)^{(\rho-1)} E_{\alpha, \beta}^{\gamma; c} \left(y \left(1 - \frac{x}{4}\right) (1-x)^2\right) dx \\ &= \left(\frac{2}{3}\right)^{2(\rho+j)} \frac{\Gamma(\rho+j)}{\Gamma(\gamma)} {}_3\psi_3 \left[\begin{array}{c} (\rho, 1), (c, 1), (\gamma, 1); \\ (\beta, \alpha), (2\rho+j, 1), (c, 1); \end{array} |y; p \right]. \quad (15) \end{aligned}$$

Proof. Let S_1 be the left-hand side of (15). Now applying the series representation of the Mittag-Leffler function (6) to the integrand of (14) and by interchanging the order of integration and summation under the given condition in Theorem (1), we

have

$$\begin{aligned}
S_1 &= \int_0^1 x^{(\rho+j-1)} (1-x)^{(2\rho-1)} \left(1 - \frac{x}{3}\right)^{(2(\rho+j)-1)} \left(1 - \frac{x}{4}\right)^{(\rho-1)} \\
&\times E_{\alpha,\beta}^{\gamma;c} \left(y \left(1 - \frac{x}{4}\right) (1-x)^2\right) dx \\
&= \int_0^1 x^{(\rho+j-1)} (1-x)^{(2\rho-1)} \left(1 - \frac{x}{3}\right)^{(2(\rho+j)-1)} \left(1 - \frac{x}{4}\right)^{(\rho-1)} \\
&\times \sum_{n=0}^{\infty} \frac{\mathbf{B}_p(\gamma+n, c-\gamma)(c)_n}{\mathbf{B}(\gamma, c-\gamma)\Gamma(\alpha n+\beta)} \frac{\left(y \left(1 - \frac{x}{4}\right) (1-x)^2\right)^n}{n!} dx \\
&= \int_0^1 x^{(\rho+j-1)} (1-x)^{2(\rho+n)-1} \left(1 - \frac{x}{3}\right)^{(2(\rho+j)-1)} \left(1 - \frac{x}{4}\right)^{(\rho+n-1)} \\
&\times \sum_{n=0}^{\infty} \frac{\mathbf{B}_p(\gamma+n, c-\gamma)(c)_n}{\mathbf{B}(\gamma, c-\gamma)\Gamma(\alpha n+\beta)} \frac{y^n}{n!} dx.
\end{aligned}$$

Now using (6) and (14), we get

$$\begin{aligned}
S_1 &= \left(\frac{2}{3}\right)^{2(\rho+j)} \frac{\Gamma(\rho+j)\Gamma(\rho+n)}{\Gamma(2\rho+j+n)} \sum_{n=0}^{\infty} \frac{\mathbf{B}_p(\gamma+n, c-\gamma)(c)_n}{\mathbf{B}(\gamma, c-\gamma)\Gamma(\alpha n+\beta)} \frac{y^n}{n!} \\
&= \left(\frac{2}{3}\right)^{(2\rho+j)} \frac{\Gamma(\rho+j)}{\Gamma(\gamma)} \sum_{n=0}^{\infty} \frac{\mathbf{B}_p(\gamma+n, c-\gamma)\Gamma(\rho+n)\Gamma(c+n)}{\Gamma(c-\gamma)\Gamma(\alpha n+\beta)\Gamma(2\rho+j+n)} \frac{y^n}{n!},
\end{aligned}$$

which upon using (12), we get

$$S_1 = \left(\frac{2}{3}\right)^{2(\rho+j)} \frac{\Gamma(\rho+j)}{\Gamma(\gamma)} {}_3\psi_3 \left[\begin{matrix} (\rho, 1), (c, 1), (\gamma, 1) \\ (\beta, \alpha), (2\rho+j, 1), (c, 1) \end{matrix}; |y; p \right].$$

□

Theorem 2. Let $y, c, \rho, j, \alpha, \beta, \gamma \in \mathbb{C}$; with $\Re(\alpha) > \max\{0, \Re(c) - 1\}; \Re(c) > 0$. $\Re(\gamma) > 0, \Re(2\rho+j) > 0$ and $x > 0$. Then the following formula holds true:

$$\begin{aligned}
&\int_0^1 x^{(\rho-1)} (1-x)^{2(\rho+j)-1} \left(1 - \frac{x}{3}\right)^{(2\rho-1)} \left(1 - \frac{x}{4}\right)^{(\rho+j-1)} E_{\alpha,\beta}^{\gamma;c} \left(yx \left(1 - \frac{x}{3}\right)^2; p\right) dx \\
&= \left(\frac{2}{3}\right)^{2(\rho)} \frac{\Gamma(\rho+j)}{\Gamma(\gamma)} {}_3\psi_3 \left[\begin{matrix} (\rho, 1), (c, 1), (\gamma, 1); \\ (\beta, \alpha), (2\rho+j, 1), (c, 1); \end{matrix}; |\frac{4}{9}y; p \right]. \quad (16)
\end{aligned}$$

Proof. Let S_2 be the left-hand side of (16). Now applying the series representation of the Mittag-Leffler function (6) to the integrand of (14), we have

$$\begin{aligned}
S_2 &= \int_0^1 x^{(\rho-1)} (1-x)^{2(\rho+j)-1} \left(1 - \frac{x}{3}\right)^{(2\rho-1)} \left(1 - \frac{x}{4}\right)^{(\rho+j-1)} \\
&\quad \times E_{\alpha,\beta}^{\gamma;c} \left(yx \left(1 - \frac{x}{3}\right)^2; p \right) dx \\
&= \int_0^1 x^{(\rho-1)} (1-x)^{2(\rho+j)-1} \left(1 - \frac{x}{3}\right)^{(2\rho-1)} \left(1 - \frac{x}{4}\right)^{(\rho+j-1)} \\
&\quad \times \sum_{n=0}^{\infty} \frac{\mathbf{B}_p(\gamma+n, c-\gamma)(c)_n}{\mathbf{B}(\gamma, c-\gamma)\Gamma(\alpha n+\beta)} \frac{\left(y \left(1 - \frac{x}{3}\right)^2\right)^n}{n!} dx \\
&= \int_0^1 x^{(\rho+n-1)} (1-x)^{2(\rho+j)-1} \left(1 - \frac{x}{3}\right)^{2(\rho+n)-1} \left(1 - \frac{x}{4}\right)^{(\rho+j-1)} \\
&\quad \times \sum_{n=0}^{\infty} \frac{\mathbf{B}_p(\gamma+n, c-\gamma)(c)_n}{\mathbf{B}(\gamma, c-\gamma)\Gamma(\alpha n+\beta)} \frac{y^n}{n!} dx.
\end{aligned}$$

Now using (6) and (14), we get

$$\begin{aligned}
S_2 &= \left(\frac{2}{3}\right)^{2(\rho+n)} \frac{\Gamma(\rho+j)\Gamma(\rho+n)}{\Gamma(2\rho+j+n)} \sum_{n=0}^{\infty} \frac{\mathbf{B}_p(\gamma+n, c-\gamma)(c)_n}{\mathbf{B}(\gamma, c-\gamma)\Gamma(\alpha n+\beta)} \frac{y^n}{n!} \\
&= \left(\frac{2}{3}\right)^{2\rho} \frac{\Gamma(\rho+j)}{\Gamma(\gamma)} \sum_{n=0}^{\infty} \frac{\mathbf{B}_p(\gamma+n, c-\gamma)\Gamma(\rho+n)\Gamma(c+n)}{\Gamma(c-\gamma)\Gamma(\alpha n+\beta)\Gamma(2\rho+j+n)} \frac{\left(\frac{4}{9}y\right)^n}{n!},
\end{aligned}$$

which upon using (12), we get

$$S_2 = \left(\frac{2}{3}\right)^{2\rho} \frac{\Gamma(\rho+j)}{\Gamma(\gamma)} {}_3\psi_3 \left[\begin{array}{l} (\rho, 1), (c, 1), (\gamma, 1); \\ (\beta, \alpha), (2\rho+j, 1), (c, 1); \end{array} ; \left| \frac{4}{9}y; p \right. \right].$$

□

Theorem 3. Let $y, c, \rho, j, \alpha, \beta, \gamma \in \mathbb{C}$; with $\Re(\alpha) > \max\{0, \Re(c) - 1\}; \Re(c) > 0$, $\Re(\gamma) > 0$, $\Re(2\rho+j) > 0$ and $x > 0$. Then the following formula holds true:

$$\begin{aligned}
&\int_0^1 x^{(\rho+j-1)} (1-x)^{(2\rho-1)} \left(1 - \frac{x}{3}\right)^{(2(\rho+j)-1)} \left(1 - \frac{x}{4}\right)^{(\rho-1)} E_{\alpha,\beta}^{\gamma;q;c} \left(y \left(1 - \frac{x}{4}\right) (1-x)^2; p \right) dx \\
&= \left(\frac{2}{3}\right)^{2(\rho+j)} \frac{\Gamma(\rho+j)}{\Gamma(\gamma)} {}_3\psi_3 \left[\begin{array}{l} (\rho, 1), (c, q), (\gamma, 1); \\ (\beta, \alpha), (2\rho+j, 1), (c, 1); \end{array} ; |y; p \right]
\end{aligned} \tag{17}$$

Proof. Let S_3 be the left-hand side of (17). Now applying the series representation of the Mittag-Leffler function (6) to the integrand of (14) and by interchanging the order of integration and summation under the given condition in Theorem (1), we

have

$$\begin{aligned}
S_3 &= \int_0^1 x^{(\rho+j-1)} (1-x)^{(2\rho-1)} \left(1-\frac{x}{3}\right)^{(2(\rho+j)-1)} \left(1-\frac{x}{4}\right)^{(\rho-1)} \\
&\times E_{\alpha,\beta}^{\gamma,q;c} \left(y \left(1-\frac{x}{4}\right) (1-x)^2; p\right) dx \\
&= \int_0^1 x^{(\rho+j-1)} (1-x)^{(2\rho-1)} \left(1-\frac{x}{3}\right)^{(2(\rho+j)-1)} \left(1-\frac{x}{4}\right)^{(\rho-1)} \\
&\times \sum_{n=0}^{\infty} \frac{\mathbf{B}_p(\gamma+nq, c-\gamma)(c)_{nq}}{\mathbf{B}(\gamma, c-\gamma)\Gamma(\alpha n+\beta)} \frac{\left(y \left(1-\frac{x}{4}\right) (1-x)^2\right)^n}{n!} dx \\
&= \int_0^1 x^{(\rho+j-1)} (1-x)^{2(\rho+n)-1} \left(1-\frac{x}{3}\right)^{(2(\rho+j)-1)} \left(1-\frac{x}{4}\right)^{(\rho+n-1)} \\
&\times \sum_{n=0}^{\infty} \frac{\mathbf{B}_p(\gamma+nq, c-\gamma)(c)_{nq}}{\mathbf{B}(\gamma, c-\gamma)\Gamma(\alpha n+\beta)} \frac{y^n}{n!} dx.
\end{aligned}$$

Now using (6) and (14), we get

$$\begin{aligned}
S_1 &= \left(\frac{2}{3}\right)^{2(\rho+j)} \frac{\Gamma(\rho+j)\Gamma(\rho+n)}{\Gamma(2\rho+j+n)} \sum_{n=0}^{\infty} \frac{\mathbf{B}_p(\gamma+nq, c-\gamma)(c)_{nq}}{\mathbf{B}(\gamma, c-\gamma)\Gamma(\alpha n+\beta)} \frac{y^n}{n!} \\
&= \left(\frac{2}{3}\right)^{(2\rho+j)} \frac{\Gamma(\rho+j)}{\Gamma(\gamma)} \sum_{n=0}^{\infty} \frac{\mathbf{B}_p(\gamma+nq, c-\gamma)\Gamma(\rho+n)\Gamma(c+nq)}{\Gamma(c-\gamma)\Gamma(\alpha n+\beta)\Gamma(2\rho+j+n)} \frac{y^n}{n!},
\end{aligned}$$

which upon using (12), we get

$$S_3 = \left(\frac{2}{3}\right)^{2(\rho+j)} \frac{\Gamma(\rho+j)}{\Gamma(\gamma)} {}_3\psi_3 \left[\begin{array}{c} (\rho, 1), (c, 1), (\gamma, 1); \\ (\beta, \alpha), (2\rho+j, 1), (c, 1); \end{array} ; |y; p \right].$$

□

Theorem 4. Let $\rho, j, \alpha, \beta, \gamma \in \mathbb{C}$; with $R(\alpha) > 0, \Re(c) > \Re(\gamma) > 0, \Re(2\rho+j) > 0$ and $x > 0$. Then the following formula holds true:

$$\begin{aligned}
&\int_0^1 x^{(\rho-1)} (1-x)^{2(\rho+j)-1} \left(1-\frac{x}{3}\right)^{(2\rho-1)} \left(1-\frac{x}{4}\right)^{(\rho+j-1)} E_{\alpha,\beta}^{\gamma,q;c} \left(yx \left(1-\frac{x}{3}\right)^2; p\right) dx \\
&= \left(\frac{2}{3}\right)^{2(\rho)} \frac{\Gamma(\rho+j)}{\Gamma(\rho)} {}_3\psi_3 \left[\begin{array}{c} (\rho, 1), (c, q), (\gamma, 1); \\ (\beta, \alpha), (2\rho+j, 1), (c, 1); \end{array} ; | \frac{4}{9}y; p \right]. \tag{18}
\end{aligned}$$

Proof. Let S_4 be the left-hand side of (18). Now applying the series representation of the Mittag-Leffler function (6) to the integrand of (14), we have

$$\begin{aligned}
S_4 &= \int_0^1 x^{(\rho-1)} (1-x)^{2(\rho+j)-1} \left(1 - \frac{x}{3}\right)^{(2\rho-1)} \left(1 - \frac{x}{4}\right)^{(\rho+j-1)} \\
&\quad \times E_{\alpha,\beta}^{\gamma,q;c} \left(yx \left(1 - \frac{x}{3}\right)^2; p\right) dx \\
&= \int_0^1 x^{(\rho-1)} (1-x)^{2(\rho+j)-1} \left(1 - \frac{x}{3}\right)^{(2\rho-1)} \left(1 - \frac{x}{4}\right)^{(\rho+j-1)} \\
&\quad \times \sum_{n=0}^{\infty} \frac{\mathbf{B}_p(\gamma+nq, c-\gamma)(c)_{nq}}{\mathbf{B}(\gamma, c-\gamma)\Gamma(\alpha n+\beta)} \frac{\left(y \left(1 - \frac{x}{3}\right)^2\right)^n}{n!} dx \\
&= \int_0^1 x^{(\rho+n-1)} (1-x)^{2(\rho+j)-1} \left(1 - \frac{x}{3}\right)^{2(\rho+n)-1} \left(1 - \frac{x}{4}\right)^{(\rho+j-1)} \\
&\quad \times \sum_{n=0}^{\infty} \frac{\mathbf{B}_p(\gamma+nq, c-\gamma)(c)_n}{\mathbf{B}(\gamma, c-\gamma)\Gamma(\alpha n+\beta)} \frac{y^n}{n!} dx
\end{aligned}$$

Now using (13) and (14), we get

$$\begin{aligned}
S_4 &= \left(\frac{2}{3}\right)^{2(\rho+n)} \frac{\Gamma(\rho+j)\Gamma(\rho+n)}{\Gamma(2\rho+j+n)} \sum_{n=0}^{\infty} \frac{\mathbf{B}_p(\gamma+nq, c-\gamma)(c)_{nq}}{\mathbf{B}(\gamma, c-\gamma)\Gamma(\alpha n+\beta)} \frac{y^n}{n!} \\
&= \left(\frac{2}{3}\right)^{2\rho} \frac{\Gamma(\rho+j)}{\Gamma(\gamma)} \sum_{n=0}^{\infty} \frac{\mathbf{B}_p(\gamma+nq, c-\gamma)\Gamma(\rho+n)\Gamma(c+nq)}{\Gamma(c-\gamma)\Gamma(\alpha n+\beta)\Gamma(2\rho+j+n)} \frac{\left(\frac{4}{9}y\right)^n}{n!},
\end{aligned}$$

which upon using (12), we get

$$S_4 = \left(\frac{2}{3}\right)^{2\rho} \frac{\Gamma(\rho+j)}{\Gamma(\gamma)} {}_3\psi_3 \left[\begin{array}{l} (\rho, 1), (c, q), (\gamma, 1); \\ (\beta, \alpha), (2\rho + j, 1), (c, 1); \\ | \frac{4}{9}y; p \end{array} \right].$$

□

3. PARTICULAR CASES

In this section, we obtained some particular cases of Theorems 1-4. Setting $p = 0$ in (6) and (13), then we get the following well known Mittag-Leffler functions.

$$E_{\alpha,\beta}^{\gamma,c}(z) = \sum_{n=0}^{\infty} \frac{\mathbf{B}_p(\gamma+n, c-\gamma)(c)_n}{\mathbf{B}(\gamma, c-\gamma)\Gamma(\alpha n+\beta)} \frac{z^n}{n!}, \quad (19)$$

and

$$E_{\alpha,\beta}^{\gamma,q;c}(z) = \sum_{n=0}^{\infty} \frac{\mathbf{B}_p(\gamma+nq, c-\gamma)(c)_{nq}}{\mathbf{B}(\gamma, c-\gamma)\Gamma(\alpha n+\beta)} \frac{z^n}{n!}. \quad (20)$$

Corollary 3.1. Let $y, c, \rho, j, \alpha, \beta, \gamma \in \mathbb{C}$; with $\Re(\alpha) > \max\{0, \Re(c) - 1\}; \Re(c) > 0$. $\Re(\gamma) > 0, \Re(2\rho + j) > 0$ and $x > 0$. Then the following formula holds true:

$$\begin{aligned} & \int_0^1 x^{(\rho+j-1)} (1-x)^{(2\rho-1)} \left(1-\frac{x}{3}\right)^{(2(\rho+j)-1)} \left(1-\frac{x}{4}\right)^{(\rho-1)} E_{\alpha,\beta}^{\gamma} \left(y \left(1-\frac{x}{4}\right) (1-x)^2\right) dx \\ &= \left(\frac{2}{3}\right)^{2(\rho+j)} \frac{\Gamma(\rho+j)}{\Gamma(\gamma)} {}_2\psi_2 \left[\begin{array}{c} (\rho, 1), (\gamma, 1); \\ (\beta, \alpha), (2\rho+j, 1); \end{array} ; |y| \right]. \end{aligned} \quad (21)$$

Corollary 3.2. Let $y, c, \rho, j, \alpha, \beta, \gamma \in \mathbb{C}$; with $\Re(\alpha) > \max\{0, \Re(c) - 1\}; \Re(c) > 0$. $\Re(\gamma) > 0, \Re(2\rho + j) > 0$ and $x > 0$. Then the following formula holds true:

$$\begin{aligned} & \int_0^1 x^{(\rho-1)} (1-x)^{2(\rho+j)-1} \left(1-\frac{x}{3}\right)^{(2\rho-1)} \left(1-\frac{x}{4}\right)^{(\rho+j-1)} E_{\alpha,\beta}^{\gamma} \left(yx \left(1-\frac{x}{3}\right)^2\right) dx \\ &= \left(\frac{2}{3}\right)^{2(\rho)} \frac{\Gamma(\rho+j)}{\Gamma(\gamma)} {}_2\psi_2 \left[\begin{array}{c} (\rho, 1), (\gamma, 1); \\ (\beta, \alpha), (2\rho+j, 1); \end{array} ; |\frac{4}{9}y| \right]. \end{aligned} \quad (22)$$

Corollary 3.3. Let $y, c, \rho, j, \alpha, \beta, \gamma \in \mathbb{C}$; with $\Re(\alpha) > \max\{0, \Re(c) - 1\}; \Re(c) > 0$. $\Re(\gamma) > 0, \Re(2\rho + j) > 0$ and $x > 0$. Then the following formula holds true:

$$\begin{aligned} & \int_0^1 x^{(\rho+j-1)} (1-x)^{(2\rho-1)} \left(1-\frac{x}{3}\right)^{(2(\rho+j)-1)} \left(1-\frac{x}{4}\right)^{(\rho-1)} E_{\alpha,\beta}^{\gamma,q} \left(y \left(1-\frac{x}{4}\right) (1-x)^2\right) dx \\ &= \left(\frac{2}{3}\right)^{2(\rho+j)} \frac{\Gamma(\rho+j)}{\Gamma(\gamma)} {}_2\psi_2 \left[\begin{array}{c} (\rho, 1), (\gamma, q); \\ (\beta, \alpha), (2\rho+j, 1); \end{array} ; |y| \right] \end{aligned} \quad (23)$$

Corollary 3.4. Let $\rho, j, \alpha, \beta, \gamma \in \mathbb{C}$; with $\Re(\alpha) > 0, \Re(c) > \Re(\gamma) > 0, \Re(2\rho+j) > 0$ and $x > 0$. Then the following formula holds true:

$$\begin{aligned} & \int_0^1 x^{(\rho-1)} (1-x)^{2(\rho+j)-1} \left(1-\frac{x}{3}\right)^{(2\rho-1)} \left(1-\frac{x}{4}\right)^{(\rho+j-1)} E_{\alpha,\beta}^{\gamma,q} \left(yx \left(1-\frac{x}{3}\right)^2\right) dx \\ &= \left(\frac{2}{3}\right)^{2(\rho)} \frac{\Gamma(\rho+j)}{\Gamma(\rho)} {}_2\psi_2 \left[\begin{array}{c} (\rho, 1), (\gamma, q); \\ (\beta, \alpha), (2\rho+j, 1); \end{array} ; |\frac{4}{9}y| \right]. \end{aligned} \quad (24)$$

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